

Non-linear Structures by Three-dimensional Dynamic Analysis

E.G. Ladopoulos
Interpaper Research Organization
8, Dimaki Str.
Athens, GR - 106 72, Greece
eladopoulos@interpaper.org

Abstract

An innovative and groundbreaking dynamic model is introduced and investigated for the solution of the three-dimensional dynamic analysis problem of a non-linear (non-symmetrical) structure, subjected under seismic forces. Such dynamic analysis problem is reduced to the solution of a system of ordinary differential equations of the second kind and the above system is numerically evaluated by using a special kind of finite elements and by solving the corresponding eigenvalues-eigenvectors problem. Finally, an application of dynamic analysis is given to the determination of the eigenvalues and eigenvectors of a 15-floor building consisting of reinforced concrete and subjected to an horizontal seismic vibration.

Key Word and Phrases

3-D Dynamic Analysis, Seismic Forces, Structural Analysis, Ordinary Differential Equations, Non-Linear Structures, Finite Elements, Multistory Frame Structures. Eigenvalues-Eigenvectors.

1. Introduction

A big variety of engineering and applied mechanics problems, have the time dimension as very important factor to be considered. So, such typical problems of engineering applications are wave transmission in fluids, transient heat conduction and dynamic behavior of structures. The latter category of applied mechanics problems, is probably the most important for practical applications, and of leading interest for dynamic analysis theory.

Consequently, because of the seismic motion of the soil, a corresponding motion is created to the frame structure. Such a seismic motion affects in many different ways the various structures. The first problem therefore to be determined is the action of an earthquake over a structure. As a second problem, the reaction of the structure over the above earthquake has to be determined. The above problems are finally reduced to the calculation of the motion for each time, for every point of the structure under study.

In addition, the study of a structural analysis problem needs a series of several steps. First of all the geometry of the structure should be extensively described and then there is a necessity for the description of the solid and section properties of the members. Also, the dynamic load conditions for which the structure needs to be analyzed, should be described.

Over the past years, E.G.Ladopoulos introduced and studied several linear [1] - [6] and non-linear singular integro-differential equations methods [7] - [10] for the solution of many important problems of structural analysis and fluid mechanics. On the other hand, an ordinary differential equations analysis is proposed by the present investigation, for the solution of some generalized dynamic analysis problems for non-linear (non-symmetrical) structures. Such problems are reduced to the solution of a system of differential equations of the second kind and such a system is finally numerically solved by using a kind of finite elements.

According to the finite element method, the continuum is divided into a finite number of elements and its behavior is specified by a finite number of parameters. The solution further of the complete system as an assembly of its elements results basically the same rules as those applicable to standard discrete problems. During the past years, many studies have been published on the application of finite elements for the numerical evaluation of problems in structural analysis and many other fields of engineering applications. Among the survey investigations on finite-element methods, are mentioned those by J.T.Oden [11] - [13], E.L.Wilson [14], K.J.Bathe and E.L.Wilson

et al. [15]-[17], O.C.Zienkiewicz [18], [19], O.C.Zienkiewicz et al. [20]-[23], H.Tottenham and C.A.Brebbia [24], J.H.Argyris et al. [25],[26], C.C.Fu [27], R.D.Kreig and S.W.Key [28], T.Belytschko, R.L.Chiapetta and H.D.Bartel [29], G.L.Goudreau and R.L.Taylor [30], T.K.Hellen [31], I.Fried [32], [33], M.Ziamal [34], L.S.D.Morley [35], J.C.Nactegaal, D.M.Parks and J.R.Rice [36], J.L.Batoz, A.Chattopadhyay and G.Dhatt [37] and T.Matsui and O.Matsuoka [38].

By the present investigation a new dynamic model is proposed for the solution of the three-dimensional structural analysis problem of a non-linear (non-symmetrical) structure under seismic forces. Such problem is reduced to the solution of a system of ordinary differential equations of the second kind and such a system is numerically evaluated by using finite elements and solving the corresponding eigenvalues-eigenvectors problem.

So, the proposed finite element method will be much smaller in degree of freedom size than commercial software, because classical linear stiffness matrices of 3-dimensional beam element have six degrees of freedom per node. The new model is therefore only applied for buildings and thus some special beam elements are used, which differs from the usual 3-D elements for structures.

An application is finally given, to the determination of the eigenvalues and eigenvectors of a 15-floor building consisting of reinforced concrete and subjected under horizontal seismic vibration. For the calculation of the above eigenvalues and eigenvectors, the Jacobi transformation method of a real symmetric matrix is used. This method consists of a sequence of orthogonal similarity transformations, known as the Jacobi rotations.

2. Three-Dimensional Non-linear Dynamic Analysis

Consider the multistory frame structure of Figure 1, under the following assumptions:

- I. The masses of this structure are concentrated in the levels of plates.
- II. The material of columns performs Hooke's law of elasticity.
- III. Simple harmonic seismic forces are applied to the frame structure.

If the above mentioned assumptions are fulfilled, then the multistory frame structure is reduced to the solution of a system of n masses $m_1, m_2, \dots, m_i, \dots, m_n$ which is concentrated in the levels of girders and connected together and with the ground, by the elastic joints $k_1, k_2, \dots, k_i, \dots, k_n$.

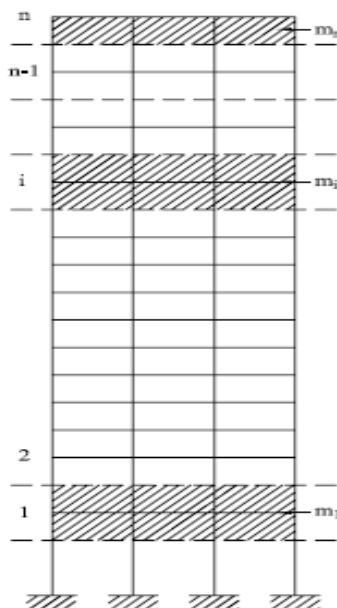


Fig.1 Multistory frame structure subjected to seismic forces

Beyond the above, consider by w_{0t} the deflection of the foundation from its initial position at time t , and w_{it} ($i=1,2,\dots,n$) the corresponding deflection of mass m_{it} ($i=1,2,\dots,n$), at the same time. Also, let by u_{it} the displacement of mass m_{it} at this time t , too (Figure 2).

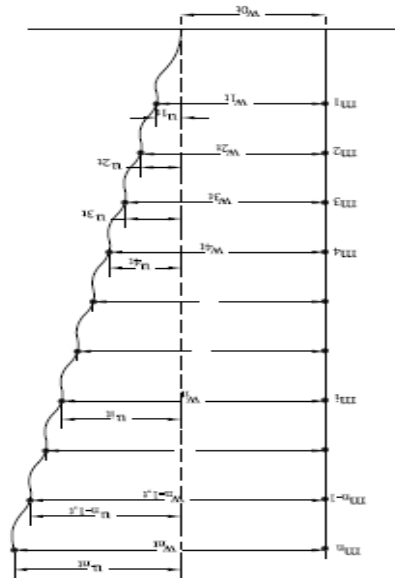


Fig. 2 The multistory frame structure is reduced to a system of masses m_i concentrated in the levels of girders

Consequently, from Figure 2 it is well seen, that following relation between the deflections and the displacements is valid:

$$w_{it} = w_{0t} + u_{it} \quad , \quad i = 1, 2, \dots, n \quad (2.1)$$

Moreover, equations of motion for the masses m_i , $i = 1, 2, \dots, n$ are considered. So, in the mass m_i , of the i -floor are effected the elastic force of restoring R_{it} , the force of damping A_{it} and the force of inertia $m_i \ddot{w}_{it}$. (Figure 3)

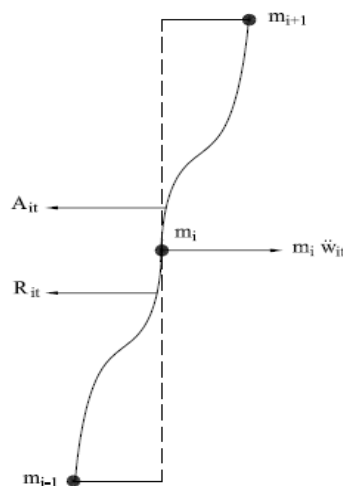


Fig. 3 Dynamic model. of multistory frame structure

and by replacing $\ddot{\mathbf{w}}_t = \ddot{\mathbf{u}}_t + \delta\ddot{w}_{0t}$ takes the form:

$$\mathbf{M}\ddot{\mathbf{u}}_t + \mathbf{C}\dot{\mathbf{u}}_t + \mathbf{K}\mathbf{u}_t = -(\mathbf{M} \cdot \delta)\ddot{w}_{0t} \quad (2.9)$$

where:

$$\mathbf{M} = \begin{bmatrix} m_1 & & & \\ & m_2 & & \\ & & \ddots & \\ & & & m_n \end{bmatrix}, \mathbf{C} = \begin{bmatrix} c_{11} & c_{12} & \dots & c_{1n} \\ c_{21} & c_{22} & \dots & c_{2n} \\ \dots & \dots & \dots & \dots \\ c_{n1} & c_{n2} & \dots & c_{nn} \end{bmatrix}, \mathbf{K} = \begin{bmatrix} k_{11} & k_{12} & \dots & k_{1n} \\ k_{21} & k_{22} & \dots & k_{2n} \\ \dots & \dots & \dots & \dots \\ k_{n1} & k_{n2} & \dots & k_{nn} \end{bmatrix} \quad (2.10)$$

with \mathbf{M} , \mathbf{C} , and \mathbf{K} the mass, damping and stiffness matrices, correspondingly.

The stiffness will be determined by following next method. Hence, consider that the displacements u_{it} are due to the external loads $P_{it} = -R_{it}$. Then because of (2.3) we obtain: (Fig. 4)

$$P_{it} = k_{i1}u_{1t} + k_{i2}u_{2t} + \dots + k_{ii}u_{it} + \dots + k_{ij}u_{jt} + \dots + k_{in}u_{nt} \quad (2.11)$$

So, from (2.11) follows that if by the use of girders all floors of the multistory frame structure are fixed, then k_{ii} will be the proper external force for the unity displacement $u_{it} = 1$ of the floor i and k_{ij} will be the corresponding reaction in the girder of the floor i for the unity displacement $u_j = 1$ of the floor j . (Fig. 4)

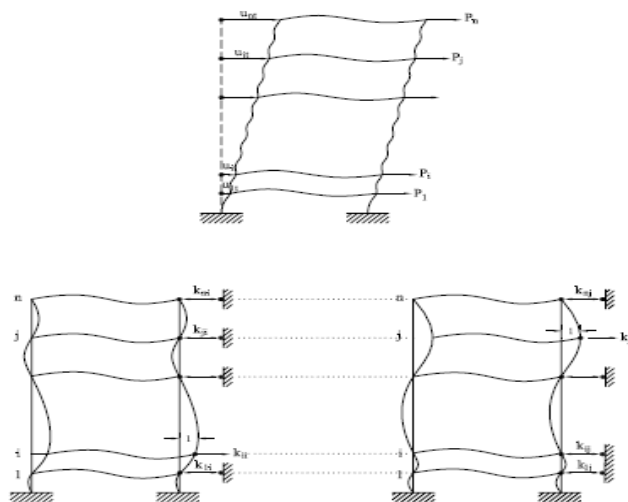


Fig. 4 Calculation of the stiffness matrix.

Consequently, by the former unity displacement of the floor i are calculated the stiffnesses $k_{\lambda i}$ ($\lambda = 1, 2, \dots, n$) of the column i of the stiffness matrix \mathbf{K} , and by the latter by the same way are calculated the stiffnesses $k_{\lambda j}$ ($j = 1, 2, \dots, n$) of column j . By successive therefore unity displacement of the total number of floors of the multistory frame structure are calculated all the columns of the stiffness matrix. On the other hand, according to Betti's reciprocal theorem the equalities $k_{ij} = k_{ji}$ are equal, and thus the stiffness matrix is always symmetric.

3. Natural Vibration and Eigenvalues for Dynamic Analysis

In order the system of ordinary differential equations (2.9) to be solved, there is a necessity for the study of the natural vibration, according to which some special solutions are calculated, which are independent from the external seismic excitation. By the above special solutions are simultaneously determined all the dynamic characteristics of the mechanical system and thus the final solution of the problem is obtained.

Hence, in order to determine the above special solutions, then the following assumptions should be considered:

- I. The damping of the system is equal to zero ($\mathbf{c} = 0$).
- II. The vibration is due to known displacements u_{i0} and velocities \dot{u}_{i0} when beginning the calculation of the time ($t = 0$).

If the above mentioned assumptions are fulfilled, then (2.9) takes the following form:

$$\mathbf{M}\ddot{\mathbf{u}}_t + \mathbf{K}\mathbf{u}_t = 0 \quad (3.1)$$

Some solutions of the following form are investigated for the system of ordinary differential equations (3.1):

$$\mathbf{u}_t = [u_{1t}, u_{2t}, \dots, u_{nt}]^T = [\varphi_1, \varphi_2, \dots, \varphi_n]^T f_t = \boldsymbol{\varphi} f_t \quad (3.2)$$

By substituting (3.2) in (3.1), follows:

$$\begin{bmatrix} m_{11} & m_{12} & \dots & m_{1n} \\ m_{21} & m_{22} & \dots & m_{2n} \\ \dots & \dots & \dots & \dots \\ m_{n1} & m_{n2} & \dots & m_{nn} \end{bmatrix} \begin{bmatrix} \varphi_1 \\ \varphi_2 \\ \dots \\ \varphi_n \end{bmatrix} \ddot{f}_t + \begin{bmatrix} k_{11} & k_{12} & \dots & k_{1n} \\ k_{21} & k_{22} & \dots & k_{2n} \\ \dots & \dots & \dots & \dots \\ k_{n1} & k_{n2} & \dots & k_{nn} \end{bmatrix} \begin{bmatrix} \varphi_1 \\ \varphi_2 \\ \dots \\ \varphi_n \end{bmatrix} f_t = \begin{bmatrix} 0 \\ 0 \\ \dots \\ 0 \end{bmatrix} \quad (3.3)$$

and the i - equation of the above system of differential equations has the following form:

$$\left(\sum_{j=1}^n m_{ij} \varphi_j \right) \ddot{f}_t + \left(\sum_{j=1}^n k_{ij} \varphi_j \right) f_t = 0 \quad (3.4)$$

or further the form:

E.G. Ladopoulos

$$\frac{\sum_{j=1}^n k_{ij} \varphi_j}{\sum_{j=1}^n m_{ij} \varphi_j} = -\frac{\ddot{f}_t}{f_t} = \omega^2 \quad (3.5)$$

where ω denotes the cyclic frequency.

Thus, from (3.5) follows:

$$\sum_{j=1}^n k_{ij} \varphi_j = \omega^2 \sum_{j=1}^n m_{ij} \varphi_j, \quad i = 1, 2, \dots, n \quad (3.6)$$

which in matrix form can be also written as:

$$\mathbf{K}\boldsymbol{\varphi} = \omega^2 \mathbf{M}\boldsymbol{\varphi} \quad (3.7)$$

or of the form:

$$[\mathbf{K} - \omega^2 \mathbf{M}]\boldsymbol{\varphi} = 0 \quad (3.8)$$

where (3.8) corresponds to a generalized eigenvalue - eigenvectors problem.

The generalized eigenvalue problem (3.8) may be further written as following:

$$\begin{bmatrix} k_{11} - \omega^2 m_{11} & k_{12} - \omega^2 m_{12} & \dots & k_{1n} - \omega^2 m_{1n} \\ k_{21} - \omega^2 m_{21} & k_{22} - \omega^2 m_{22} & \dots & k_{2n} - \omega^2 m_{2n} \\ \dots & \dots & \dots & \dots \\ k_{n1} - \omega^2 m_{n1} & k_{n2} - \omega^2 m_{n2} & \dots & k_{nn} - \omega^2 m_{nn} \end{bmatrix} \begin{bmatrix} \varphi_1 \\ \varphi_2 \\ \dots \\ \varphi_n \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ \dots \\ 0 \end{bmatrix} \quad (3.9)$$

which is an homogeneous system of n linear equations with unknowns the displacements φ_i and the parameter ω^2 . In order this system to have non-zero solutions, except the obvious solution $\varphi_i = 0$, its determinant should be zero:

$$\begin{vmatrix} k_{11} - \omega^2 m_{11} & k_{12} - \omega^2 m_{12} & \dots & k_{1n} - \omega^2 m_{1n} \\ k_{21} - \omega^2 m_{21} & k_{22} - \omega^2 m_{22} & \dots & k_{2n} - \omega^2 m_{2n} \\ \dots & \dots & \dots & \dots \\ k_{n1} - \omega^2 m_{n1} & k_{n2} - \omega^2 m_{n2} & \dots & k_{nn} - \omega^2 m_{nn} \end{vmatrix} = 0 \quad (3.10)$$

This determinant will have an expansion of an algebraic equation of n - degree, with unknowns the ω^2 , from the solutions of which are calculated the following n - eigenvalues ω_j^2 :

E.G. Ladopoulos

$$\omega^2 = \omega_j^2, \quad j = 1, 2, \dots, n \quad (3.11)$$

So, in each cyclic frequency ω_j corresponds a natural period:

$$T_j = \frac{2\pi}{\omega_j} \quad (3.12)$$

Beyond the above, the homogeneous system (3.9) has n - solutions which are the eigenvectors φ_j :

$$\boldsymbol{\varphi}_j = [\varphi_{1j}, \varphi_{2j}, \dots, \varphi_{nj}]^T, \quad j = 1, 2, \dots, n \quad (3.13)$$

in correspondence to the n eigenvalues ω_j^2 .

The above homogeneous system is solved by an arbitrary selection of one component of the eigenvector φ_j (like the first one $\varphi_{1j} = c$ or the last one $\varphi_{nj} = c$) and hence the system (3.9) becomes non-homogeneous and from the solution of any subsystem with $(n-1)$ equations are further calculated the rest $(n-1)$ components of the φ_j .

Thus, the system (3.1) has n -independent solutions of the form (3.2):

$$\mathbf{u}_{jt} = [\varphi_{1j}, \varphi_{2j}, \dots, \varphi_{nj}]^T f_{jt} = \boldsymbol{\varphi}_j f_{jt}, \quad j = 1, 2, \dots, n \quad (3.14)$$

which corresponds to simple harmonic vibrations with cyclic frequencies ω_j (or periods $T_j = 2\pi/\omega_j$) and these are the natural vibrations of the mechanical system.

The above eigenvalues can be classified with increasing order of cyclic frequencies. Consequently, the eigenvalue which belongs to the lowest cyclic frequency ω_1 and correspondingly to the biggest natural period (fundamental natural period) is known as the fundamental eigenvalue.

4. Dynamic Analysis Application of a Multistory Frame Structure

Let us consider a 15-floor building consisting of reinforced concrete and which has six similar frames and hence a level of symmetry parallel to the level of the frames (Fig. 5). The dynamic characteristics of this building will be calculated for an horizontal seismic vibration.

For the above frame structure following data are valid:

$L = 5.0m$, $L_1 = 10.0m$, $L_2 = 7.0m$, $h_1 = 4.0m$, $h_2 = h_3 = \dots = h_{10} = 3.0m$, moment of inertia of all the columns $I = 0.02m^4$ (70×70), moment of inertia of T - beams $J = 0.022m^4$ (beams 30×70 ,

E.G. Ladopoulos

Because of symmetry the stiffness matrix of the building follows by a simple sum of the stiffness matrices \mathbf{K}_0 of the six frames. Hence, Fig. 6 shows the direct calculation of the horizontal stiffness matrix of a frame which corresponds to three freedoms of movement of its floors.

For example by enforcing the unity condition:

$$u_1 = 1, \quad u_2 = u_3 = \dots = u_{15} = 0 \quad (4.2)$$

and applying any static method, by using one kind of finite-elements [Falter: 39], then k_{11} is calculated as the requested external force for the unity displacement $u_1 = 1.0m$ of the first floor and k_{i1} , $i = 1, 2, \dots, 15$ are calculated as the reactions in the node i of the floor i , $i = 1, 2, \dots, 15$ because of the unity displacement $u_1 = 1.0m$ of the first floor.

So, the results of the static calculations are:

$$\begin{aligned} k_{11} &= 695,471kN \\ k_{21} &= -458,457kN \\ k_{31} &= 22,107kN \\ k_{41} &= 5599kN \\ k_{51} &= 131kN \\ k_{61} &= -219kN \\ k_{71} &= -14kN \\ k_{81} &= 7kN \\ k_{91} &= 1.5kN \\ k_{101} &= 0.4kN \\ k_{111} &= 0.5kN \\ k_{121} &= 0.4kN \\ k_{131} &= -4kN \\ k_{141} &= -11kN \\ k_{151} &= 96kN \end{aligned} \quad (4.3)$$

By the same way and by enforcing the unity conditions $u_2 = 1, u_3 = 1, \dots, u_{15} = 1$ are calculated the values of the coefficients of stiffness $(k_{12}, k_{22}, \dots, k_{152}), (k_{13}, k_{23}, \dots, k_{153}), \dots, (k_{110}, k_{210}, \dots, k_{1515})$.

Hence, the stiffness matrix of the building takes the following form:

$\mathbf{K} =$

695,471	-458,457	22,107	5599	131	-219	-14	7	1.5	0.4	0.5	0.4	-4	-11	96
-458,457	860,111	-457,404	22,020	5458	134	-215	-14	-6	1.3	0.4	0.3	-3	-10	83
22,107	-457,404	860,168	-457,521	22,033	5,462	134	-215	-14	6	1.3	0.2	-3	-10	81
5599	22,020	-457,521	860,161	-457,520	22,034	5462	133	-215	-14	6	1.1	-3	-12	82
131	5458	22,033	-457,520	860,164	-457,523	22,035	5462	133	-215	-14	6	-2	-10	82
-219	134	5462	22,034	-457,523	860,165	-457,525	22,036	5462	133	-215	-14	3	-10	83
-14	-215	134	5462	22,035	-457,525	860,167	-457,527	22,036	5462	133	-215	-17	-4.2	84
7	-14	-215	133	5462	22,036	-457,527	860,168	-457,529	22,037	5462	133	-219	-25	91
1.5	6	-14	-215	133	5462	22,036	-457,529	860,169	-457,560	22,038	5463	127	-230	80
0.4	1.3	6	-14	-215	133	5462	22,037	-457,560	860,166	-457,530	22,040	5465	113	-132
0.5	0.4	1.3	6	-14	-215	133	5462	22,038	-457,530	860,166	-457,531	22,084	5541	-49
0.4	0.3	0.2	1.1	6	-14	-215	133	5463	22,040	-457,531	860,146	-457,474	22397	5304
-4	-3	-3	-3	-2	3	-17	-219	127	5465	22,084	-457,474	859,260	-458,714	28,075
-11	-10	-10	-12	-10	-10	-4.2	-25	-230	113	5541	22,397	-458,714	844,912	-420,702
96	83	81	82	82	83	84	91	80	-132	-49	5304	28,075	-420,702	394,538

(4.4)

From eqn. (4.4) can be seen that the stiffness matrix is symmetrical.

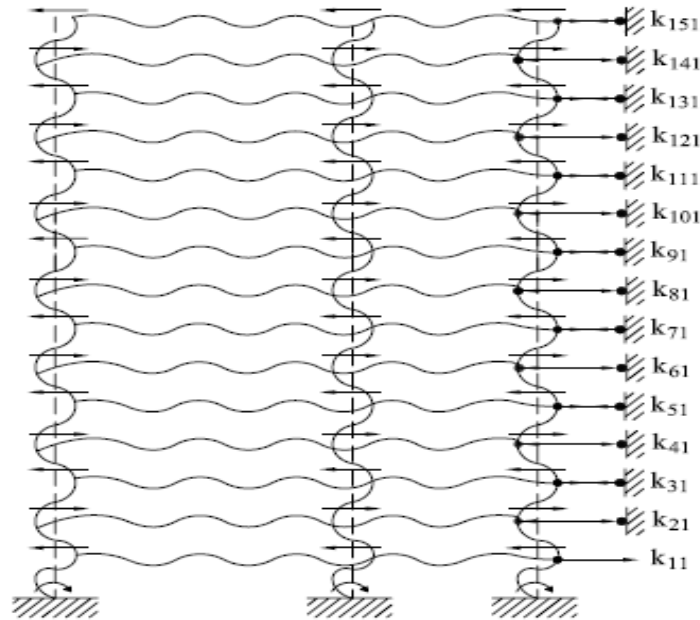


Fig. 6 Direct calculation of the stiffness matrix.

Beyond the above, by using the expressions for the mass and the stiffness matrices given by eqs (4.1) and (4.4), then the characteristic equation (3.10) of the eigenvalue-eigenvectors problem (3.8) reduces to the following form:

$$|\mathbf{K} - m\omega^2\mathbf{M}_0| =$$

E.G. Ladopoulos

695,471-λ	-458,457	22,107	5599	131	-219	-14	7	1.5	0.4	0.5	0.4	-4	-11	96
-458,457	860,111-λ	-457,404	22,020	5458	134	-215	-14	-6	1.3	0.4	0.3	-3	-10	83
22,107	-457,404	860,168-λ	-457,521	22,033	5,462	134	-215	-14	6	1.3	0.2	-3	-10	81
5599	22,020	-457,521	860,161-λ	-457,520	22,034	5462	133	-215	-14	6	1.1	-3	-12	82
131	5458	22,033	-457,520	860,164-λ	-457,523	22,035	5462	133	-215	-14	6	-2	-10	82
-219	134	5462	22,034	-457,523	860,165-λ	-457,525	22,036	5462	133	-215	-14	3	-10	83
-14	-215	134	5462	22,035	-457,525	860,167-λ	-457,527	22,036	5462	133	-215	-17	-4.2	84
7	-14	-215	133	5462	22,036	-457,527	860,168-λ	-457,529	22,037	5462	133	-219	-25	91
1.5	6	-14	-215	133	5462	22,036	-457,529	860,169-λ	-457,560	22,038	5463	127	-230	80
0.4	1.3	6	-14	-215	133	5462	22,037	-457,560	860,166-λ	-457,530	22,040	5465	113	-132
0.5	0.4	1.3	6	-14	-215	133	5462	22,038	-457,530	860,166-λ	-457,531	22,084	5541	-49
0.4	0.3	0.2	1.1	6	-14	-215	133	5463	22,040	-457,531	860,146-λ	-457,474	22397	5304
-4	-3	-3	-3	-2	3	-17	-219	127	5465	22,084	-457,474	859,260-λ	-458,714	28,075
-11	-10	-10	-12	-10	-10	-4.2	-25	-230	113	5541	22,397	-458,714	844,912-λ	-420,702
96	83	81	82	82	83	84	91	80	-132	-49	5304	28,075	-420,702	394,538-λ

(4.5)

where:

$$\lambda = m\omega^2 \quad (4.6)$$

Moreover, by using the Jacobi transformations method of a real symmetrical matrix [40], then all eigenvalues are computed. The above method consists of a sequence of orthogonal similarity transformations. Each transformation known as a Jacobi rotation, is only a plane rotation designed to annihilate one of the off-diagonal matrix elements. The successive transformations undo previously set zeros, but the off-diagonal elements are getting smaller and smaller, until the matrix finally becomes diagonal to the computer precision.

Thus, the eigenvalues which are computed are:

$$\begin{aligned} \lambda_1 &= 3046 \text{ KN/m}^2, \lambda_2 = 27,872 \text{ KN/m}^2, \lambda_3 = 79,188 \text{ KN/m}^2, \\ \lambda_4 &= 158,143 \text{ KN/m}^2, \lambda_5 = 266,042 \text{ KN/m}^2, \lambda_6 = 402,322 \text{ KN/m}^2, \\ \lambda_7 &= 564,379 \text{ KN/m}^2, \lambda_8 = 746,338 \text{ KN/m}^2, \lambda_9 = 939,734 \text{ KN/m}^2, \\ \lambda_{10} &= 1,134,090 \text{ KN/m}^2, \lambda_{11} = 1,318,720 \text{ KN/m}^2, \lambda_{12} = 1,483,780 \text{ KN/m}^2 \\ \lambda_{13} &= 1,620,980 \text{ KN/m}^2, \lambda_{14} = 1,723,810 \text{ KN/m}^2, \lambda_{15} = 1,787,480 \text{ KN/m}^2 \end{aligned}$$

For the eigenvalues we have:

$$\begin{aligned} \omega_1^2 &= \frac{\lambda_1}{m} = \frac{3046}{780} \Rightarrow \omega_1 = 1.98 \text{ rad/sec}, T_1 = 3.18 \text{ sec} \\ \omega_2^2 &= \frac{\lambda_2}{m} = \frac{27,872}{780} \Rightarrow \omega_2 = 5.98 \text{ rad/sec}, T_2 = 1.05 \text{ sec} \\ \omega_3^2 &= \frac{\lambda_3}{m} = \frac{79,188}{780} \Rightarrow \omega_3 = 10.08 \text{ rad/sec}, T_3 = 0.62 \text{ sec} \\ \omega_4^2 &= \frac{\lambda_4}{m} = \frac{158,143}{780} \Rightarrow \omega_4 = 14.24 \text{ rad/sec}, T_4 = 0.44 \text{ sec} \\ \omega_5^2 &= \frac{\lambda_5}{m} = \frac{266,042}{780} \Rightarrow \omega_5 = 18.47 \text{ rad/sec}, T_5 = 0.34 \text{ sec} \\ \omega_6^2 &= \frac{\lambda_6}{m} = \frac{402,322}{780} \Rightarrow \omega_6 = 22.71 \text{ rad/sec}, T_6 = 0.28 \text{ sec} \\ \omega_7^2 &= \frac{\lambda_7}{m} = \frac{564,379}{780} \Rightarrow \omega_7 = 26.90 \text{ rad/sec}, T_7 = 0.23 \text{ sec} \\ \omega_8^2 &= \frac{\lambda_8}{m} = \frac{746,338}{780} \Rightarrow \omega_8 = 30.93 \text{ rad/sec}, T_8 = 0.20 \text{ sec} \end{aligned}$$

E.G. Ladopoulos

$$\begin{aligned}\omega_9^2 &= \frac{\lambda_9}{m} = \frac{939,734}{780} \Rightarrow \omega_9 = 34.71 \text{ rad/sec}, T_9 = 0.18 \text{ sec} \\ \omega_{10}^2 &= \frac{\lambda_{10}}{m} = \frac{1,134,090}{780} \Rightarrow \omega_{10} = 38.13 \text{ rad/sec}, T_{10} = 0.16 \text{ sec} \\ \omega_{11}^2 &= \frac{\lambda_{11}}{m} = \frac{1,318,720}{780} \Rightarrow \omega_{11} = 41.12 \text{ rad/sec}, T_{11} = 0.15 \text{ sec} \\ \omega_{12}^2 &= \frac{\lambda_{12}}{m} = \frac{1,483,780}{780} \Rightarrow \omega_{12} = 43.62 \text{ rad/sec}, T_{12} = 0.14 \text{ sec} \\ \omega_{13}^2 &= \frac{\lambda_{13}}{m} = \frac{1,620,980}{780} \Rightarrow \omega_{13} = 45.59 \text{ rad/sec}, T_{13} = 0.14 \text{ sec} \\ \omega_{14}^2 &= \frac{\lambda_{14}}{m} = \frac{1,723,810}{780} \Rightarrow \omega_{14} = 47.01 \text{ rad/sec}, T_{14} = 0.13 \text{ sec} \\ \omega_{15}^2 &= \frac{\lambda_{15}}{m} = \frac{1,787,480}{780} \Rightarrow \omega_{15} = 47.87 \text{ rad/sec}, T_{15} = 0.13 \text{ sec}\end{aligned}$$

The eigenvectors are further calculated from the homogeneous system of linear equations (3.9), by using the Jacobi transformations method [40]:

$$(\mathbf{K} - \lambda_i \mathbf{M}_0) \boldsymbol{\varphi}_i =$$

695,471 - λ_i	-458,457	22,107	5599	131	-219	-14	7	1.5	0.4	0.5	0.4	-4	-11	96
-458,457	860,111 - λ_i	-457,404	22,020	5458	134	-215	-14	-6	1.3	0.4	0.3	-3	-10	83
22,107	-457,404	860,168 - λ_i	-457,521	22,033	5,462	134	-215	-14	6	1.3	0.2	-3	-10	81
5599	22,020	-457,521	860,161 - λ_i	-457,520	22,034	5462	133	-215	-14	6	1.1	-3	-12	82
131	5458	22,033	-457,520	860,164 - λ_i	-457,523	22,035	5462	133	-215	-14	6	-2	-10	82
-219	134	5462	22,034	-457,523	860,165 - λ_i	-457,525	22,036	5462	133	-215	-14	3	-10	83
-14	-215	134	5462	22,035	-457,525	860,167 - λ_i	-457,527	22,036	5462	133	-215	-17	-4.2	84
7	-14	-215	133	5462	22,036	-457,527	860,168 - λ_i	-457,529	22,037	5462	133	-219	-25	91
1.5	6	-14	-215	133	5462	22,036	-457,529	860,169 - λ_i	-457,560	22,038	5463	127	-230	80
0.4	1.3	6	-14	-215	133	5462	22,037	-457,560	860,166 - λ_i	-457,530	22,040	5465	113	-132
0.5	0.4	1.3	6	-14	-215	133	5462	22,038	-457,530	860,166 - λ_i	-457,531	22,084	5541	-49
0.4	0.3	0.2	1.1	6	-14	-215	133	5463	22,040	-457,531	860,146 - λ_i	-457,474	22397	5304
-4	-3	-3	-3	-2	3	-17	-219	127	5465	22,084	-457,474	859,260 - λ_i	-458,714	28,075
-11	-10	-10	-12	-10	-10	-4.2	-25	-230	113	5541	22,397	-458,714	844,912 - λ_i	-420,702
96	83	81	82	82	83	84	91	80	-132	-49	5304	28,075	-420,702	394,538 - λ_i

$$\begin{matrix} \begin{bmatrix} \varphi_{1i} \\ \varphi_{2i} \\ \varphi_{3i} \\ \varphi_{4i} \\ \varphi_{5i} \\ \varphi_{6i} \\ \varphi_{7i} \\ \varphi_{8i} \\ \varphi_{9i} \\ \varphi_{10i} \\ \varphi_{11i} \\ \varphi_{12i} \\ \varphi_{13i} \\ \varphi_{14i} \\ \varphi_{15i} \end{bmatrix} \\ \times \end{matrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix} \quad (4.7)$$

for $\varphi_{15i} = 1.0, i = 1, 2, \dots, 15$

Consequently, by solving the above homogeneous system, then next eigenvalues are obtained:

E.G. Ladopoulos

$$\begin{aligned}
 \Psi_1 &= \begin{bmatrix} \varphi_{11} \\ \varphi_{21} \\ \varphi_{31} \\ \varphi_{41} \\ \varphi_{51} \\ \varphi_{61} \\ \varphi_{71} \\ \varphi_{81} \\ \varphi_{91} \\ \varphi_{101} \\ \varphi_{111} \\ \varphi_{121} \\ \varphi_{131} \\ \varphi_{141} \\ \varphi_{151} \end{bmatrix} = \begin{bmatrix} 0.153 \\ 0.254 \\ 0.349 \\ 0.441 \\ 0.528 \\ 0.611 \\ 0.688 \\ 0.758 \\ 0.822 \\ 0.877 \\ 0.925 \\ 0.964 \\ 0.994 \\ 1.010 \\ 1.000 \end{bmatrix} &
 \Psi_2 &= \begin{bmatrix} \varphi_{12} \\ \varphi_{22} \\ \varphi_{32} \\ \varphi_{42} \\ \varphi_{52} \\ \varphi_{62} \\ \varphi_{72} \\ \varphi_{82} \\ \varphi_{92} \\ \varphi_{102} \\ \varphi_{112} \\ \varphi_{122} \\ \varphi_{132} \\ \varphi_{142} \\ \varphi_{152} \end{bmatrix} = \begin{bmatrix} -0.442 \\ -0.698 \\ -0.887 \\ -0.997 \\ -1.022 \\ -0.960 \\ -0.816 \\ -0.602 \\ -0.336 \\ -0.041 \\ 0.257 \\ 0.534 \\ 0.756 \\ 0.928 \\ 1.000 \end{bmatrix} &
 \Psi_3 &= \begin{bmatrix} \varphi_{13} \\ \varphi_{23} \\ \varphi_{33} \\ \varphi_{43} \\ \varphi_{53} \\ \varphi_{63} \\ \varphi_{73} \\ \varphi_{83} \\ \varphi_{93} \\ \varphi_{103} \\ \varphi_{113} \\ \varphi_{123} \\ \varphi_{133} \\ \varphi_{143} \\ \varphi_{153} \end{bmatrix} = \begin{bmatrix} 0.678 \\ 0.971 \\ 1.024 \\ 0.831 \\ 0.441 \\ -0.053 \\ -0.535 \\ -0.889 \\ -1.033 \\ -0.932 \\ -0.610 \\ -0.143 \\ 0.358 \\ 0.772 \\ 1.000 \end{bmatrix} \\
 \\
 \Psi_4 &= \begin{bmatrix} \varphi_{14} \\ \varphi_{24} \\ \varphi_{34} \\ \varphi_{44} \\ \varphi_{54} \\ \varphi_{64} \\ \varphi_{74} \\ \varphi_{84} \\ \varphi_{94} \\ \varphi_{104} \\ \varphi_{114} \\ \varphi_{124} \\ \varphi_{134} \\ \varphi_{144} \\ \varphi_{154} \end{bmatrix} = \begin{bmatrix} -0.856 \\ -1.040 \\ -0.735 \\ -0.088 \\ 0.600 \\ 1.013 \\ 0.962 \\ 0.470 \\ -0.237 \\ -0.836 \\ -1.050 \\ -0.783 \\ -0.157 \\ 0.542 \\ 1.000 \end{bmatrix} &
 \Psi_5 &= \begin{bmatrix} \varphi_{15} \\ \varphi_{25} \\ \varphi_{35} \\ \varphi_{45} \\ \varphi_{55} \\ \varphi_{65} \\ \varphi_{75} \\ \varphi_{85} \\ \varphi_{95} \\ \varphi_{105} \\ \varphi_{115} \\ \varphi_{125} \\ \varphi_{135} \\ \varphi_{145} \\ \varphi_{155} \end{bmatrix} = \begin{bmatrix} 0.994 \\ 0.930 \\ 0.164 \\ -0.729 \\ -1.080 \\ -0.628 \\ 0.291 \\ 0.994 \\ 0.958 \\ 0.211 \\ -0.694 \\ -1.082 \\ -0.666 \\ 0.248 \\ 1.000 \end{bmatrix} &
 \Psi_6 &= \begin{bmatrix} \varphi_{16} \\ \varphi_{26} \\ \varphi_{36} \\ \varphi_{46} \\ \varphi_{56} \\ \varphi_{66} \\ \varphi_{76} \\ \varphi_{86} \\ \varphi_{96} \\ \varphi_{106} \\ \varphi_{116} \\ \varphi_{126} \\ \varphi_{136} \\ \varphi_{146} \\ \varphi_{156} \end{bmatrix} = \begin{bmatrix} -1.107 \\ -0.669 \\ 0.502 \\ 1.136 \\ 0.545 \\ -0.634 \\ -1.129 \\ -0.405 \\ 0.756 \\ 1.101 \\ 0.258 \\ -0.864 \\ -1.053 \\ -0.100 \\ 1.000 \end{bmatrix}
 \end{aligned}$$

(4.8)

E.G. Ladopoulos

$$\Phi_7 = \begin{bmatrix} \varphi_{17} \\ \varphi_{27} \\ \varphi_{37} \\ \varphi_{47} \\ \varphi_{57} \\ \varphi_{67} \\ \varphi_{77} \\ \varphi_{87} \\ \varphi_{97} \\ \varphi_{107} \\ \varphi_{117} \\ \varphi_{127} \\ \varphi_{137} \\ \varphi_{147} \\ \varphi_{157} \end{bmatrix} = \begin{bmatrix} 1.209 \\ 0.284 \\ -1.062 \\ -0.869 \\ 0.587 \\ 1.190 \\ 0.063 \\ -1.155 \\ -0.693 \\ 0.777 \\ 1.118 \\ -0.166 \\ -1.209 \\ -0.487 \\ 1.000 \end{bmatrix}$$

$$\Phi_8 = \begin{bmatrix} \varphi_{18} \\ \varphi_{28} \\ \varphi_{38} \\ \varphi_{48} \\ \varphi_{58} \\ \varphi_{68} \\ \varphi_{78} \\ \varphi_{88} \\ \varphi_{98} \\ \varphi_{108} \\ \varphi_{118} \\ \varphi_{128} \\ \varphi_{138} \\ \varphi_{148} \\ \varphi_{158} \end{bmatrix} = \begin{bmatrix} -1.303 \\ 0.209 \\ 1.342 \\ -0.005 \\ -1.341 \\ -0.194 \\ 1.313 \\ 0.389 \\ -1.255 \\ -0.575 \\ 1.169 \\ 0.749 \\ -1.059 \\ -0.898 \\ 1.000 \end{bmatrix}$$

$$\Phi_9 = \begin{bmatrix} \varphi_{19} \\ \varphi_{29} \\ \varphi_{39} \\ \varphi_{49} \\ \varphi_{59} \\ \varphi_{69} \\ \varphi_{79} \\ \varphi_{89} \\ \varphi_{99} \\ \varphi_{109} \\ \varphi_{119} \\ \varphi_{129} \\ \varphi_{139} \\ \varphi_{149} \\ \varphi_{159} \end{bmatrix} = \begin{bmatrix} 1.392 \\ -0.785 \\ -1.195 \\ 1.087 \\ 0.915 \\ -1.322 \\ -0.575 \\ 1.470 \\ 0.198 \\ -1.520 \\ 0.193 \\ 1.471 \\ -0.572 \\ -1.315 \\ 1.000 \end{bmatrix}$$

$$\Phi_{10} = \begin{bmatrix} \varphi_{110} \\ \varphi_{210} \\ \varphi_{310} \\ \varphi_{410} \\ \varphi_{510} \\ \varphi_{610} \\ \varphi_{710} \\ \varphi_{810} \\ \varphi_{910} \\ \varphi_{1010} \\ \varphi_{1110} \\ \varphi_{1210} \\ \varphi_{1310} \\ \varphi_{1410} \\ \varphi_{1510} \end{bmatrix} = \begin{bmatrix} -1.471 \\ 1.412 \\ 0.552 \\ -1.768 \\ 0.604 \\ 1.373 \\ -1.501 \\ -0.392 \\ 1.758 \\ -0.756 \\ -1.264 \\ 1.581 \\ 0.229 \\ -1.722 \\ 1.000 \end{bmatrix}$$

$$\Phi_{11} = \begin{bmatrix} \varphi_{111} \\ \varphi_{211} \\ \varphi_{311} \\ \varphi_{411} \\ \varphi_{511} \\ \varphi_{611} \\ \varphi_{711} \\ \varphi_{811} \\ \varphi_{911} \\ \varphi_{1011} \\ \varphi_{1111} \\ \varphi_{1211} \\ \varphi_{1311} \\ \varphi_{1411} \\ \varphi_{1511} \end{bmatrix} = \begin{bmatrix} 1.540 \\ -2.049 \\ 0.557 \\ 1.474 \\ -2.067 \\ 0.644 \\ 1.408 \\ -2.086 \\ 0.728 \\ 1.340 \\ -2.100 \\ 0.811 \\ 1.268 \\ -2.102 \\ 1.000 \end{bmatrix}$$

$$\Phi_{12} = \begin{bmatrix} \varphi_{112} \\ \varphi_{212} \\ \varphi_{312} \\ \varphi_{412} \\ \varphi_{512} \\ \varphi_{612} \\ \varphi_{712} \\ \varphi_{812} \\ \varphi_{912} \\ \varphi_{1012} \\ \varphi_{1112} \\ \varphi_{1212} \\ \varphi_{1312} \\ \varphi_{1412} \\ \varphi_{1512} \end{bmatrix} = \begin{bmatrix} -1.596 \\ 2.651 \\ -1.991 \\ 0.048 \\ 1.927 \\ -2.656 \\ 1.668 \\ 0.399 \\ -2.208 \\ 2.589 \\ -1.295 \\ -0.836 \\ 2.425 \\ -2.439 \\ 1.000 \end{bmatrix}$$

$$\begin{array}{l}
 \Phi_{13} = \begin{bmatrix} \varphi_{113} \\ \varphi_{213} \\ \varphi_{313} \\ \varphi_{413} \\ \varphi_{513} \\ \varphi_{613} \\ \varphi_{713} \\ \varphi_{813} \\ \varphi_{913} \\ \varphi_{1013} \\ \varphi_{1113} \\ \varphi_{1213} \\ \varphi_{1313} \\ \varphi_{1413} \\ \varphi_{1513} \end{bmatrix} = \begin{bmatrix} 1.641 \\ -3.174 \\ 3.522 \\ -2.557 \\ 0.637 \\ 1.522 \\ -3.112 \\ 3.539 \\ -2.643 \\ 0.760 \\ 1.408 \\ -3.049 \\ 3.549 \\ -2.717 \\ 1.000 \end{bmatrix} \\
 \Phi_{14} = \begin{bmatrix} \varphi_{114} \\ \varphi_{214} \\ \varphi_{314} \\ \varphi_{414} \\ \varphi_{514} \\ \varphi_{614} \\ \varphi_{714} \\ \varphi_{814} \\ \varphi_{914} \\ \varphi_{1014} \\ \varphi_{1114} \\ \varphi_{1214} \\ \varphi_{1314} \\ \varphi_{1414} \\ \varphi_{1514} \end{bmatrix} = \begin{bmatrix} -1.673 \\ 3.581 \\ -4.883 \\ 5.362 \\ -4.993 \\ 3.668 \\ -1.783 \\ -0.405 \\ 2.524 \\ -4.215 \\ 5.192 \\ -5.290 \\ 4.491 \\ -2.926 \\ 1.000 \end{bmatrix} \\
 \Phi_{15} = \begin{bmatrix} \varphi_{115} \\ \varphi_{215} \\ \varphi_{315} \\ \varphi_{415} \\ \varphi_{515} \\ \varphi_{615} \\ \varphi_{715} \\ \varphi_{815} \\ \varphi_{915} \\ \varphi_{1015} \\ \varphi_{1115} \\ \varphi_{1215} \\ \varphi_{1315} \\ \varphi_{1415} \\ \varphi_{1515} \end{bmatrix} = \begin{bmatrix} 1.692 \\ -3.834 \\ 5.815 \\ -7.550 \\ 8.961 \\ -9.990 \\ 10.590 \\ -10.737 \\ 10.425 \\ -9.665 \\ 8.491 \\ -6.953 \\ 5.115 \\ -3.054 \\ 1.000 \end{bmatrix}
 \end{array}$$

The above eigenvalues are giving the final solution of the eigenvalue-eigenvectors problem (3.8). The big practical value of the modal analysis consists on the possibility of expressing of every dynamic behavior of a building as a linear superposition of its eigenvalues.

5. Conclusions

In general, there is very little known about the procedure which makes earthquakes. Several forces of much different kind try continuously to change the present state of earth. These forces are external, i.e. these which get energy from sources outside the earth (like solar energy, total energy, etc.) and internal, i.e. these which their energy rises from the interior of the earth (for example gravity, the mechanical forces due to the rotation of the earth, the loss of heat, the nuclear energy, etc.).

Hence, the action of these forces leads to the accumulation of elastic stresses between the masses of the rocks. As the accumulation of these stresses is continuously effected, then the strength of the rocks is overdrawed and hence its equilibrium is disturbed and some sudden changes are created. So, a big part of the dynamic energy due to the stresses is changed to kinetic under the form of seismic wavings.

The most important problem of dynamic analysis is therefore to calculate the motion for each time and for every point of the structure. Thus, by the current research a new dynamic model has been proposed for the solution of the three-dimensional structural analysis problem of a non-linear structure subjected under seismic forces. This problem of dynamic analysis was finally reduced to the solution of a system of ordinary differential equations of the second kind.

Some kind of finite elements were used for the solution of this system of ordinary differential equations and solving the corresponding eigenvalues-eigenvectors problem. An application was finally given to the determination of the eigenvalues and eigenvectors of a 15-floor building consisting of reinforced concrete and subjected under horizontal seismic vibration. For the calculation of the above eigenvalues and eigenvectors, the Jacobi transformation method of a real symmetric matrix was used. This method consists of a sequence of orthogonal similarity transformations, known as the Jacobi rotations.

References

1. E.G.Ladopoulos, 'On a new integration rule with the Gegenbauer polynomials for singular integral equations, used in the theory of elasticity', *Ing. Archiv*, **58** (1988), 35-46.

E.G. Ladopoulos

2. E.G.Ladopoulos, 'On the numerical evaluation of the multidimensional singular integrals and integral equations used in the theory of linear viscoelasticity', *Int. J. Math. Math. Scien.*, **11** (1988), 561-574.
3. E.G.Ladopoulos, 'Singular integral operators method for three-dimensional elasto-plastic stress analysis', *Comp. Struct.*, **38** (1991), 1-8.
4. E.G.Ladopoulos, 'Singular integral operators method for two-dimensional elasto-plastic stress analysis', *Forsch. Ingen.*, **57** (1991), 152-158.
5. E.G.Ladopoulos, 'Singular integral operators method for anisotropic elastic stress analysis', *Comp. Struct.*, **48** (1993), 965-973.
6. E.G.Ladopoulos, 'Singular Integral Equations, Linear and Non-linear Theory and its Applications in Science and Engineering', Springer-Verlag, Berlin, New York, 2000.
7. E.G.Ladopoulos, 'Non-linear integro-differential equations used in orthotropic shallow spherical shell analysis', *Mech. Res. Commun.*, **18** (1991), 111-119.
8. E.G.Ladopoulos, 'Non-linear integro-differential equations in sandwich plates stress analysis', *Mech. Res. Commun.*, **21** (1994), 95-102.
9. E.G.Ladopoulos, 'Non-linear singular integral computational analysis for unsteady flow problems', *Renew. Energy*, **6** (1995), 901-906.
10. E.G.Ladopoulos, 'Non-linear singular integral representation analysis for inviscid flowfields of unsteady airfoils', *Int. J. Non-Lin. Mech.*, **32** (1997), 377-384.
11. J.T.Oden, 'A general theory of finite elements - I: Topological considerations', *Int. J. Num. Meth. Engng*, **1** (1969), 205-221.
12. J.T.Oden, 'A general theory of finite elements - II: Applications', *Int. J. Num. Meth. Engng.*, **1** (1969), 247-260.
13. J.T.Oden, 'Finite Elements of Non-linear Continua', McGraw-Hill, Berkshire, 1971.
14. E.L.Wilson, 'Solid SAP - A static analysis program for three-dimensional solid structures', *SESM Report 71-19*, Dept. Civil Engng, Univ. California, Berkeley (1971).
15. K.J.Bathe and E.L.Wilson, 'Stability and accuracy analysis of direct integration methods', *Int. J. Earth. Eng. Struct. Dynam.*, **1** (1973), 283-291.
16. K.J.Bathe, E.L.Wilson and R.H.Iding, 'NONSAP - A structural analysis program for static and dynamic response of nonlinear systems', *SESM Report 74-3*, Dept. Civil Engng, Univ. California, Berkeley (1974).
17. K.J.Bathe, E.L.Wilson and F.E.Peterson, 'SAPIV - A structural analysis program for static and dynamic response of linear systems', *Report EERC 73-11*, Dept. Civil Engng, Univ. California, Berkeley (1974).
18. O.C.Zienkiewicz, 'Constrained variational principles and penalty function methods in the finite element analysis', *Lecture Notes in Mathematics*, No.363, p.p. 207-314, Springer-Verlag, New York (1974).
19. O.C.Zienkiewicz, 'The Finite Element Method', McGraw-Hill, Berkshire, 1994.
20. O.C.Zienkiewicz and C.J.Parekh, 'Transient field problems - two and three dimensional analysis by isoparametric finite elements', *Int. J. Num. Meth. Engng.*, **2** (1970), 61-71.
21. O.C.Zienkiewicz and R.H.Lewis, 'An analysis of various time stepping schemes for initial value problems', *Int. J. Earth. Eng. Struct. Dynam.*, **1** (1973), 407-408.
22. I.Christie, D.F.Griffiths, A.R.Mitchell and O.C.Zienkiewicz, 'Finite element methods for second order equations with significant first derivatives', *Int. J. Num. Meth. Engng*, **10** (1976), 1389-1396.
23. O.C.Zienkiewicz, D.W.Kelly and P.Bettess, 'The coupling of the finite element method and boundary solution procedures', *Int. J. Num. Meth. Engng*, **11** (1977), 355-375.
24. H.Tottenham and C.A.Brebbia, 'Finite Element Techniques in Structural Mechanics', Southampton Univ. Press, Southampton, 1970.
25. J.H.Argyris, I.Fried and D.W.Scharpf, 'The TUBA family of plate elements for the matrix displacement method', *Aeronaut. J.*, **72** (1968), 701-709.
26. J.H.Argyris and G.Mareczek, 'Finite element analysis of slow incompressible viscous fluid motion', *Ing. Archiv.*, **43** (1974), 92-109.
27. C.C.Fu, 'On the stability of explicit methods for numerical integration of the equations of matrices in finite element methods', *Int. J. Num. Meth. Engng*, **4** (1972), 95-107.
28. R.D.Kreig and S.W.Key, 'Transient shock response by numerical time integration', *Int. J. Num. Meth. Engng*, **7** (1973), 273-286.
29. T.Belytschko, R.L.Chiapetta and H.D.Bartel, 'Efficient large scale non-linear transient analysis by finite elements', *Int. J. Num. Meth. Engng*, **10** (1976), 579-596.
30. G.L.Goudreau and R.L.Taylor, 'Evaluation of numerical integration methods in elastodynamics', *Comp. Meth. Appl. Mech. Engng*, **2** (1972), 69-97.

E.G. Ladopoulos

31. T.K.Hellen, 'Effective quadrature rules for quadratic solid isoparametric finite elements', *Int. J. Num. Meth. Engng*, **4** (1972), 597-600.
32. I.Fried, 'Accuracy and condition of curved (isoparametric) finite elements', *J. Sound Vibr.*, **31** (1973), 345-355.
33. I.Fried, 'Numerical integration in the finite element method', *Comp. Struct.*, **4** (1974), 921-932.
34. M.Ziamal, 'Curved elements in the finite element method', *SIAM J. Num. Anal.*, **11** (1974), 347-362.
35. L.S.D.Morley, 'On the constant moment plate bending element', *J. Strain Anal.* **6** (1971), 20-24.
36. J.C.Nagtegaal, D.M.Parks and J.R.Rice, 'On numerically accurate finite element solutions in the fully plastic range', *Comp. Meth. Appl. Mech. Engng*, **4** (1974), 153-178.
37. J.L.Batoz, A.Chattopadhyay and G.Dhatt, 'Finite element large deflection analysis of shallow shells', *Int. J. Num. Meth. Engng*, **10** (1976), 35-38.
38. T.Matsui and O.Matsuoka, 'A new finite element scheme for instability analysis of thin shells', *Int. J. Num. Meth. Engng*, **10** (1976), 145-170.
39. B.Falter, 'Statikprogramme für Personalcomputer', Werner Verlag, Düsseldorf, 1992.
40. W.H.Press, S.A.Teukolsky, W.T.Vetterling and B.P.Flannery, 'Numerical Recipes in Fortran 77: The Art of Scientific Computing', Second Edit., Cambridge Univ. Press, Cambridge, 1999.